

*Short note*

## First observation of excited states in the neutron-rich nucleus $^{138}\text{Te}$

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**Abstract.** The very neutron-rich nucleus  $^{138}\text{Te}$  produced in the spontaneous fission of  $^{248}\text{Cm}$  has been investigated using the EUROGAM II  $\gamma$ -ray multidetector array. The excited states of  $^{138}\text{Te}$  observed for the first time indicate that the region of ground state deformation beyond the doubly magic nucleus  $^{132}\text{Sn}$  has yet not been reached.

**PACS.** 23.20.Lv Gamma transitions and level energies – 25.85.Ca Spontaneous fission – 27.60.+j  $90 \leq A \leq 149$

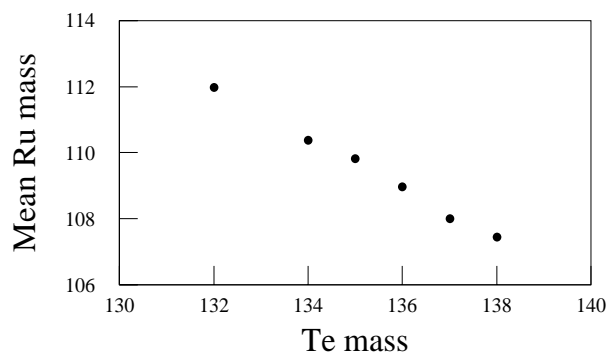
Shape evolutions in nuclei with mass numbers near 132 are strongly influenced by shell closures. An indication of the stable, axially-symmetric, spheroidal deformation in nuclei is given by the ratio of the excitation energies of the first  $4^+$  and  $2^+$  levels and it amounts to 3.33 for good rotational nuclei. On the neutron-rich side of the doubly magic nucleus  $^{132}\text{Sn}$ , the locus of the transition from single-particle to collective behaviour is unknown. For example, the four-valence particle nucleus  $^{136}\text{Te}$  displays a level scheme typical of single-particle behaviour. In order to search for the expected transition, the present work investigates the very neutron-rich nucleus  $^{138}\text{Te}$ .

The  $^{138}\text{Te}$  nuclei were obtained as secondary fragments in the spontaneous fission of  $^{248}\text{Cm}$ . The source consisted of about 5 mg of  $^{248}\text{Cm}$ , in the form of oxide, embedded uniformly in a pellet of potassium chloride. The radioactive source was placed in the centre of the EUROGAM II array [1], which consisted in this experiment of 52 escape-suppressed Ge detectors including 24 four-crystal CLOVER detectors and which was augmented by the addition of 4 low-energy photon spectrometers. A total of approximately  $2.5 \times 10^9$  threefold or higher coincidence events was collected within a period of 11 days.

The level scheme of  $^{138}\text{Te}$  was constructed by examining double gated monodimensional  $\gamma$  spectra. Information

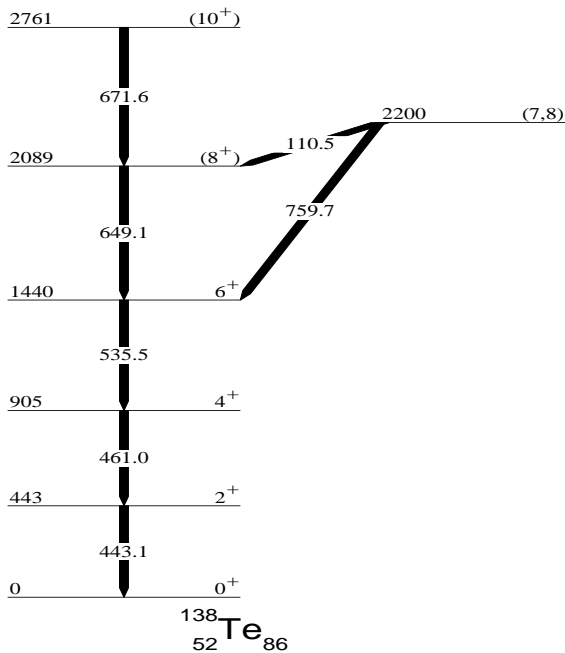
on the multiplicities of  $\gamma$  rays could not be obtained due to the fact that the yield of  $^{138}\text{Te}$  in the spontaneous fission of  $^{248}\text{Cm}$  is too weak. Since no  $\gamma$  lines in  $^{138}\text{Te}$  were known prior to the present work, the  $\gamma$  lines were identified through the coincidence relationship with the complementary Ru fission fragments. The result of applying this identification technique, which is described in detail in previous publications [2,3], is shown in Fig. 1.

The partial level scheme of  $^{138}\text{Te}$  obtained in this work is shown in Fig. 2. The ordering of the levels is based

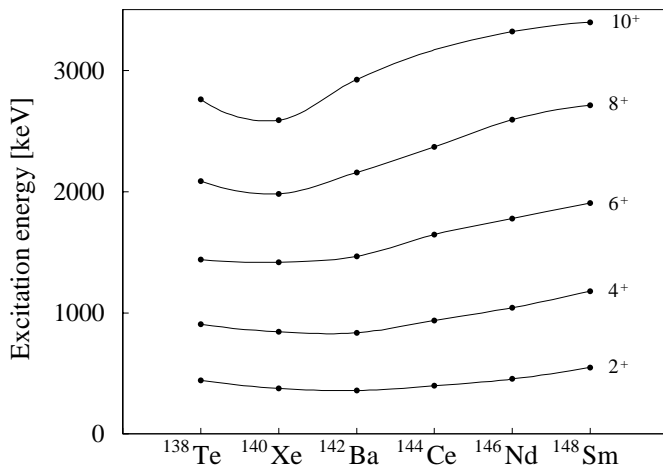


**Fig. 1.** Plot of the mean mass of Ru fragments which accompany specified Te fragments. The smooth trend in the data points is the basis for the assignment of new  $\gamma$ -rays lines to  $^{138}\text{Te}$

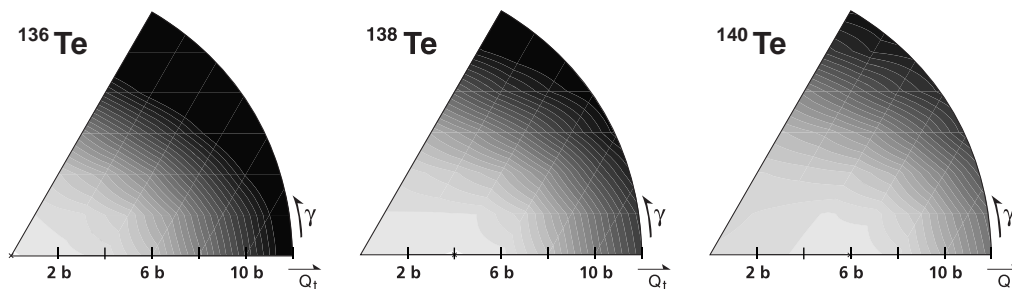
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**Fig. 2.** Level scheme of  $^{138}\text{Te}$  obtained in the present work. Level and  $\gamma$ -ray energies are given in keV



**Fig. 3.** Systematics of levels for the ground state bands of  $N = 86$  isotones. Data are taken from this work and [3–7]



**Fig. 4.** Potential energy surfaces in the  $(Q_t, \gamma)$ -plane for  $^{136,138,140}\text{Te}$  calculated via Hartree-Fock plus BCS code. The distance between contour lines is 500 keV. The minima are indicated by asterisks

on relative intensities of  $\gamma$  rays. Spin assignments of  $2^+$ ,  $4^+$  and  $6^+$  to the three first excited states are based on the systematics of the ground-state bands in neighbouring  $N=86$  isotones as can be seen in Fig. 3.

The ratio of the excitation energies of the first  $4^+$  and  $2^+$  levels,  $E(4^+)/E(2^+) = 2.04$ , indicates that  $^{138}\text{Te}$  is not a strongly deformed nucleus. A more precise analysis, namely the description of states in closed and near-closed shell nuclei, uses the plots of  $E(I^+)/E(2^+)$  versus  $E(4^+)/E(2^+)$ , called also Mallmann's plots [8]. The values for  $E(6^+)/E(2^+)$  and  $E(8^+)/E(2^+)$ , 3.26 and 4.71, respectively, are in agreement with the  $\alpha$  root of [8] and point to a vibrational description of the  $^{138}\text{Te}$  nucleus.

Concurrently, Hartree-Fock plus BCS calculations [9] based on the new SLy4 Skyrme interaction [10] and a density dependent  $\delta$  pairing [11] indicate that the  $^{138}\text{Te}$  nucleus lies just at the transition between a spherical  $^{136}\text{Te}$  and a prolate  $^{140}\text{Te}$ . Moreover, as seen in Fig. 4,  $^{138}\text{Te}$  exhibits a very  $\beta$  soft prolate minimum consistent with the present experimental data.

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